

Research Article

Sodium Adsorption, Diffusion and Coverage on Two-Dimensional GeP₃

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Abstract

In this study, we performed density functional theory calculations to investigate Na adsorption and diffusion properties on two-dimensional germanium phosphide (III) (GeP₃) monolayer. We found that single Na atom can strongly be adsorbed (2.71 eV) on the top of the valley site Ge atom and leads to a semiconductor to metal transition in GeP₃. The adsorption energy of Na decreases by increased Na concentration. With full Na concentration, the GeP₃ monolayer is subject to an increase in lattice parameter of about 0.25%. This increase rate is very small as desired in ion batteries. Moreover, Na atoms have very small diffusion barrier energy (0.25 eV) on GeP₃ monolayers and the calculated diffusion coefficient at room temperature is $1.15 \times 10^{-7} \text{ cm}^2/\text{s}$. In addition, calculated average open-circuit voltage is 1.48 eV, which is comparable with the other Na-ion anode materials. Our results show that GeP₃ monolayer can be a good candidate anode material for Na-ion batteries with a maximum capacity of 324 mAhg⁻¹.

Introduction

Increasing demand for clean energy has encouraged researchers to develop novel energy

storage devices, such as Lithium-ion (Li-ion) batteries. There are numerous experimental and theoretical studies on Li-ion batteries that focus on using two-dimensional (2D) materials as electrodes [1-6]. Several studies have revealed that graphene can store more Li atoms than graphite due to the former's increased volume-surface ratio [7-9]. However, experimental studies

show that transition metal dichalcogenides (TMDs) as MoS₂ have higher Li capacities than graphene [10]. Ersan, et al. predicted that monolayer TMDs MX₂ (M=Mo, W and X=O, S, Se, and Te) have lower activation energy barriers than graphene or silicene [11]. These lower energy barrier values result in higher diffusion coefficient for Li atoms, so that mobility of Li atoms on MX₂ monolayers can be higher than many commercial anode materials. In addition, many theoretical studies focus on lithium adsorption on 2D semiconductor materials to

investigate the electronic structure transition from semiconductor to metal [12,13].

However, limited Li resources necessitate new elements that can be used in energy storage. Since sodium (Na) and Li are in the same group in the Periodic table, they have similar properties. Hence, Na can be an excellent candidate for future large-scale energy storage applications. As mentioned above, graphene is suitable anode material for Li-ion battery applications with relatively high Li storage capacity [14]. In contrast, Malyi et al. showed that ideal graphene cannot be used as an anode for Na-ion batteries [15]. However, synthesized antimony/multilayer graphene heterostructure shows a high reversible Na storage capacity with a value of 452 mAhg⁻¹ [16]. In another study, the synthesis of antimony carbon nanocomposite alloy was reported by Qian and co-workers [17]. It was found that this material has a long-term cycling stability with 94% capacity retention over 100 cycles, and it has Na storage capacity of 610 mAhg⁻¹. By performing first-principles Density Functional Theory (DFT) simulations, Sengupta showed that 2D AlN can form a good Na-ion battery with a 385 mAhg⁻¹ Na storage capacity, 1.86 V open-circuit voltage (OCV), and 0.15 eV diffusion barrier [18]. Similar to AlN (1.49 V), BN (1.46 V) has suitable cell voltage for Na-ion batteries and these cell voltages are comparable to carbon nanotubes [19]. Theoretical studies have shown that the maximum Na storage capacity of layered MoS₂ is 146 mAhg⁻¹, which is low than its Li storage capacity [20]. Increasing of Na coverage on MoS₂ induces a phase transformation from 1H-MoS₂ to 1T-MoS₂.

Recently, a new 2D material GeP₃ has been proposed by Heine et al [21]. The bulk form of GeP₃ is metallic and has layered structure. Also, the calculated cleavage energy for monolayer

GeP₃ is 1.14 Jm⁻² [21]. Adsorption and diffusion of Li atom on this newly predicted monolayer has been investigated by Zhang et al., and Li capacity has been reported as 648 mAhg⁻¹ [22]. In addition, 2D GeP₃ is found also can be good material as a high capacity anode material for non-Lithium-ion batteries [23]. Furthermore, it has also be pointed out that this monolayer can be a good candidate for gas sensor applications such as CH₄, N₂, and H₂ molecules [24]. InP₃ monolayer, which has the same structure as GeP₃, is a promising anode material for both Li- and Na-ion batteries with a high storage capacity of 258 mAhg⁻¹ [25].

In this study, we investigated Na adsorption and diffusion on 2D GeP₃ monolayers by performing first-principles DFT calculations. Although Li and Na are in the same column of the Periodic table and accordingly share similar principles and properties, investigations have been focused mostly on Li-ion batteries instead of Na-ion batteries. Besides, Na resources are almost unlimited in the Earth's crust and ocean, so, Na-ion batteries are more economical than Li-ion batteries. Therefore, exploration of the nature of Na-ion batteries will be very important to improving energy storage applications.

Computational Details

In this study, first-principles DFT is used to investigate adsorption and diffusion properties of a Na atom on the surface of a GeP₃ monolayer. The calculations are performed using spin-polarized, plane-wave DFT with ultrasoft pseudopotentials [26] as implemented in Quantum ESPRESSO software package [27]. For the exchange-correlation functional, generalized gradient approximation (GGA) is employed by using Perdew-Burke-Ernzerhof (PBE) parametrization [28]. The van der Waals (vdW) corrections are taken into account by using DFT-

D2 [29]. For the self-consistent field potential and total energy calculations, the Brillouin Zone (BZ) is sampled using a $(12 \times 12 \times 1)$ Monkhorst-Pack special k-points mesh for monolayer GeP_3 unit-cell [30]. A plane-wave basis set with a kinetic energy cutoff of 70 Ry is used for all calculations. The lattice constants and atomic positions are optimized by using the BFGS quasi-Newton algorithm [31]. The convergence criterion for energy is chosen as 10^{-6} Ry between two successive iteration steps, and the maximum Hellmann-Feynman forces exerted on each atom is reduced to less than $0.005 \text{ Ry}/\text{\AA}$ upon ionic relaxation. To avoid interlayer interactions, two GeP_3 monolayers are separated by a distance of 20\AA vacuum length. We used the climbing image nudged elastic band (CI-NEB) method to seek the minimum energy path and saddle points between determined initial and final positions [32].

In order to determine the adsorption energy E_{ads} of Na atoms on a GeP_3 monolayer we used the following equation;

$$E_{\text{ads}} = nE_{\text{Na}} + E_{\text{GeP}_3} - E_{\text{GeP}_3+\text{Na}} \quad (1)$$

where E_{Na} and E_{GeP_3} are the total energy of the isolated Na atom and the total energy of bare GeP_3 monolayer, respectively. $E_{\text{GeP}_3+\text{Na}}$ is the total energy of the Na atom adsorbed GeP_3 monolayer.

Results and Discussion

For the sake of comparison, we first carried out the properties of a monolayer structure of GeP_3 . The calculated optimized lattice constants for the 2×2 supercell are $a=b=13.87 \text{\AA}$. We also evaluated the cohesive energy for GeP_3 monolayer by $E_{\text{coh}} = (mE_{\text{Ge}} + nE_{\text{P}} - E_{\text{GeP}_3}) / (m + n)$, where E_{GeP_3} is the total energy of the bare monolayer, and E_{Ge}

and E_{P} are the total energies of isolated free atoms of Ge and P, respectively. m and n are the numbers of Ge and P atoms in the unit cell. The calculated cohesive energy is 3.43 eV/atom . Fig. 1 depicts the atomic structure, band structure, and total and partial electronic density of states of the monolayer GeP_3 . The PBE calculation predicted that the band gap of the GeP_3 monolayer is 0.22 eV and the band character is indirect from

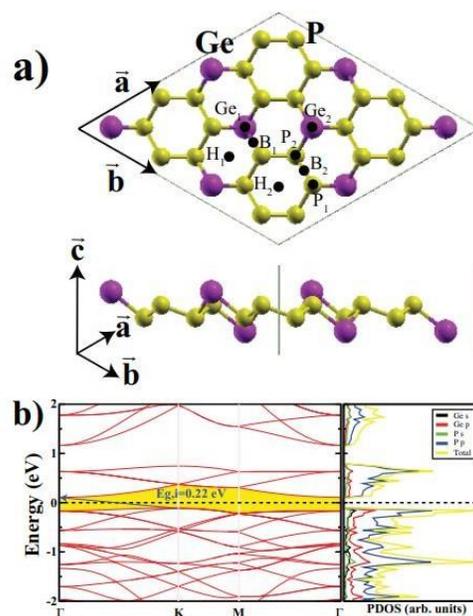


Figure 1: (Color online) a) Top and side views of the GeP_3 monolayer, eight considered adsorption sites on the GeP_3 monolayer are also shown. b) The electronic band structure and corresponding total and partial density of states of the 2×2 supercell of GeP_3 monolayer. The Fermi level is set to zero energy. Indirect band character is shown by a blue arrow.

the K (valence band maximum; VBM) to the Γ (conduction band minimum; CBM). Both the CBM and VBM are occupied dominantly by phosphorus p orbitals; the second highest

contribution comes from germanium p orbitals. These calculated values are consistent with the previous results [21,22,24].

We next investigated the adsorption of a single Na atom on the GeP₃ monolayer. We used (2×2) supercells to prevent the Na-Na interactions between the neighbor cells. With this enlargement, the corresponding chemical stoichiometry is Ge₈P₂₄Na. As illustrated in Figure 1(a), eight adsorption sites were considered to determine the most energetically favorable positions for a single Na atom bound to GeP₃. These sites are: the top of Ge atoms (Ge₁ is for valley site Ge atoms, Ge₂ is for the top of upper-level Ge atoms), top of P atoms (P₁, P₂), top of hollow sites (H₁, H₂) and top of the bridge sites (B₁, B₂). Our calculations show that the most favorable adsorption site for Na atom is Ge₁ with a E_{ads} value of 2.71 eV. The most favorable adsorption sites from best to worst are as follows: Ge₁=P₁ > H₁ = H₂ > P₂ > Ge₂. The corresponding adsorption energies are 2.71 eV > 2.65 eV > 2.60 eV > 2.02 eV. During geometrical optimization, a Na atom on a P₁ site slips towards a Ge₁ site, and a Na atom on a H₁ site also slips close to a Ge₁ site. These adsorption energies prove that there is a strong interaction between the Na atom and the GeP₃ monolayer.

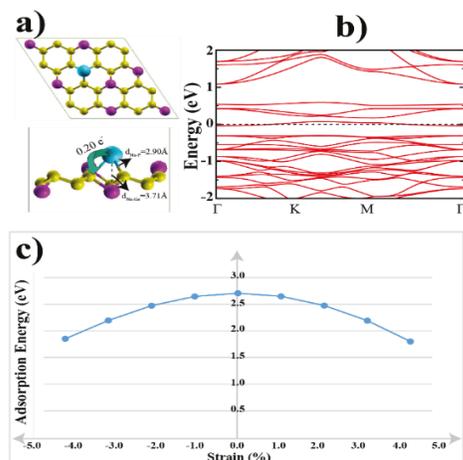


Figure 2: (Color online) a) Top and side views of the single Na atom adsorbed GeP₃ monolayer with bond distances. The arrow shows the electron transfers from Na atom to the surface. b) The electronic band structure of the Na adsorbed GeP₃ system. c) Adsorption energies of a single Na atom as the function of applied biaxial tensile/compressive strain percentage.

Figure 2(a) and (b) illustrate the atomic structure of the Na atom adsorbed GeP₃ monolayer with the bond distances and electronic band dispersion after the geometric optimization. Löwdin analysis indicates that 0.20 electrons transferred from Na atom to the GeP₃ monolayer while Na atom on Ge₁ site. This amount of charge transfer results in the shifting of the Fermi energy level of the GeP₃ and semiconducting-metal transition occurs and eventually, the new structure becomes a non-magnetic metal which has a lower resistivity and higher electronic conductivity than bare GeP₃ (see Figure 2(b)). The lattice constants are decreased only 0.14% by single Na atom adsorption and this extremely small volume change during Na intercalation may solve the volume expansion issue in Na-ion batteries. To see how the applied tensile or compressive strain affect the binding energy of Na atom, we applied biaxial in-plane strains to the monolayer GeP₃. From Figure 2(c), we can see that the adsorption energy decreases quickly from 2.71 eV to 1.80 eV with tensile or compressive strains. This can imply that Na atom adsorption energy is sensitive to the external strain.

Even though low electrical resistivity is important to determine the rate performance of anode materials, it is not sufficient by itself. The adsorbed atom diffusion performance should also be known. Therefore, the diffusion of Na atom on GeP₃ monolayer is also examined in this study. To seek the minimum energy path and saddle points,

we use the climbing image nudged elastic band (CI-NEB) method [32]. For these calculations,

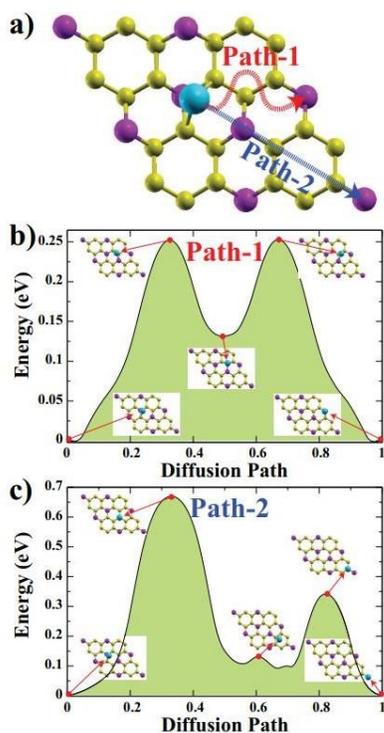


Figure 3: (Color online) a) Top view of the considered diffusion paths for Na (Path-1 (red arrow), Path-2 (blue arrow)), b) Calculated energy profiles along the Path-1 and c) Path-2 for one Na

atom on the GeP₃ monolayer. Insets depict the atomic structures for the transition states in Path-1 and Path-2.

we determined two linear different paths to determine the activation barrier energy of the Na atom. Besides, the selected initial and final images are fully optimized before the CI-NEB calculation. All diffusion paths divided by 23 equal intervals, at the end of the CI-NEB calculations the paths that Na atom follows can see in Figure 3(a) with red and blue arrows.

For both Path-1 and Path-2, there are two energy minima, one of them is absolute energy minimum where the Na atom is adsorbed on Ge₁ site, and the other one is a local minimum corresponding to the Na atom close to H₂ site. The two maxima occur while Na atom is passing from the B₁ sites for the Path-1. Due to Ge₂ site is the least favorable adsorption site among the all considered sites, high energy barrier happens when the Na atom is on this site (see inset of Figure 2(c)). The calculated activation energy barriers for these two paths for Na atom are 0.25 eV and 0.66 eV for Path-1 and Path-2, respectively. These energy barriers are comparable with Na atom on InP₃ monolayer which is obtained by Liu, et al. [25] From these CI-NEB calculations we can easily say that Na atom prefers to diffuse on the line of Path-1.

The diffusion coefficient (D) at the room temperature of the Na atom on GeP₃ can be calculated using Arrhenius equation by the following equation:

$$D \simeq a^2 \nu \exp\left(\frac{-E_a}{k_B T}\right) \quad (2)$$

where a is the lattice constant of GeP₃ monolayer ($a=6.935\text{\AA}$), $\nu=1 \times 10^{11}$ Hz (jumping frequency), and $k_B T=0.026$ eV. The calculated diffusion coefficient along the Path-1 is $1.15 \times 10^{-7} \text{ cm}^2/\text{s}$ and this D value indicates that Na atom can easily move on the GeP₃ which is important for the electrochemical reactions in ion batteries.

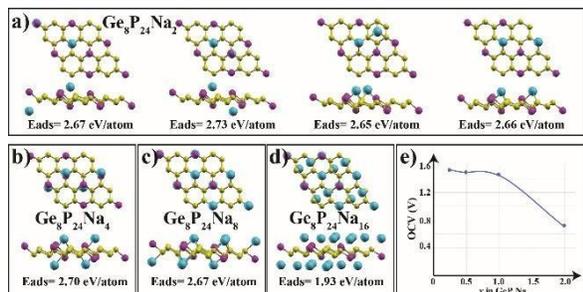


Figure 4: (Color online) a) Top and side views of the optimized atomic structures of Ge₈P₂₄Na₂ (GeP₃Na_{0.25}), b) Ge₈P₂₄Na₄ (GeP₃Na_{0.5}), c) Ge₈P₂₄Na₈ (GeP₃Na), and d) Ge₈P₂₄Na₁₆ (GeP₃Na₂), e) The electrode potential profiles of Na interactions on GeP₃.

As is known, the performance of the Na(Li)-ion batteries is directly related to the storage capacity. For this reason, we further examined more Na atoms adsorbed on both sides of the GeP₃ monolayer to specify the highest theoretical Na atom-specific capacity. For this investigation, we first considered two individual Na atoms in two main scenarios. In the first scenario, the Na atoms are adsorbed to different sides of the same GeP₃ monolayer. In the other case, both Na atoms are adsorbed to the same side. Both cases are illustrated in Figure 4(a). Our calculations show that the adsorption energy of the Na atoms on opposite sides is higher by 60 meV/atom than that of the Na atoms on one side. This implies that two Na atoms prefer to bind GeP₃ separately on both sides to counteract the possible structural distortion in GeP₃. Figure 4(b), (c), and (d) show the atomic structures and adsorption energies for the increased number of Na atoms on the GeP₃ monolayer. After full coverage of the Na atoms on the Ge₁ sites (see Figure 4(c)), we placed eight more Na atoms on P₂ sites on top of the Na atoms level. After geometrical optimization, these additional Na atoms penetrated and remained between the lower level Na atoms without

creating Na clusters (see Figure 4(d)). Thus all Na atoms were positioned on the same level for each side of the GeP₃ monolayer.

The maximum theoretical capacities (C) are estimated by the following equation:

$$C_M = \frac{xF}{M_{\text{GeP}_3}} \quad (3)$$

where M_{GeP_3} is the atomic mass of the GeP₃ monolayer, x is the highest number of adsorbed Na atoms, and F is the Faraday constant. The calculated maximum Na storage capacity for the GeP₃ monolayer is 324 mAhg⁻¹ which is higher than that of InP₃ monolayer (258 mAhg⁻¹) [25], because of the relatively light atomic weight of the Ge atom. In addition, this Na storage capacity on GeP₃ is larger than those on many transition metal tellurides (MTe₂; M=Co, Fe, Mn, Sc, Ti : 22-130 mAhg⁻¹) [33], amorphous (Glassy) carbon (173 mAhg⁻¹) [34] and MoS₂ (146 mAhg⁻¹) [35]. From these results, we can say that if the GeP₃ monolayer structure is used as an anode material for Na ion batteries, it could show high power density performance. Also, the full coverage of Na atoms on GeP₃ results in just a 0.25% lattice constant enhancement for a 2x2 supercell of GeP₃, which is a negligible change in the volume of the cell.

Disregarding entropic and enthalpic energy contributions, the average open-circuit voltage (V_{OCV}) can be calculated from total-energy differences with the following equation:

$$V_{OCV} = \frac{x E_{\text{Na}} + E_{\text{GeP}_3} - E_{\text{GeP}_3+x\text{Na}}}{xye} \quad (4)$$

where $E_{\text{GeP}_3+x\text{Na}}$ and E_{GeP_3} are the total energies of GeP_3 with and without Na intercalation. E_{Na} is the cohesive energy of metal Na; x is the number of adsorbed Na atom on the GeP_3 and y is the electrical charge of Na ions in the electrolyte (we assume that $y = 1$ in this equation). Figure 4(e) shows the calculated average open-circuit voltage profile of Na on GeP_3 . As the Na atom concentration increases from 0.25 to 2.0, the electrode potential (V) tends to decrease.

However, we should note that the biggest decrease occurs between GeP_3Na_2 to GeP_3Na , while for other concentrations the potential alteration is only 0.05 V. The calculated average V_{OCV} for $\text{GeP}_3\text{Na}_{0.25}$ to GeP_3Na is 1.48 V, which is similar to the commercial anode material of TiO_2 (1.50V) [36].

Conclusion

In conclusion, we systematically investigated the adsorption and diffusion of sodium on the GeP_3 monolayer using first-principles density functional calculations. We also studied the external strain effects on a single Na atom

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adsorbed on a GeP_3 monolayer, and found that adsorption energy of Na atom on GeP_3 is reduced by tensile or compressive strain. The semiconducting GeP_3 monolayer turns to metal upon adsorption of a single Na atom, which leads to an increase of electrical conductivity of the material. With full Na concentration, the GeP_3 monolayer undergoes only a 0.25% lattice parameter increase and the theoretical Na storage capacity is 324 mAhg^{-1} . Besides, the calculated average V_{OCV} from $\text{GeP}_3\text{Na}_{0.25}$ to GeP_3Na is 1.48 V, and is comparable to the electrode potential of the commercial anode materials. Our results indicate that the GeP_3 monolayer is a promising anode material for Na-ion battery applications similar to other newly predicted two-dimensional materials [37,42].

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